

Chem 140a

Lecture Notes #8: Class #10

**Recombination-Generation Statistics (R-G)
(Shockley-Read-Hall statistics)**

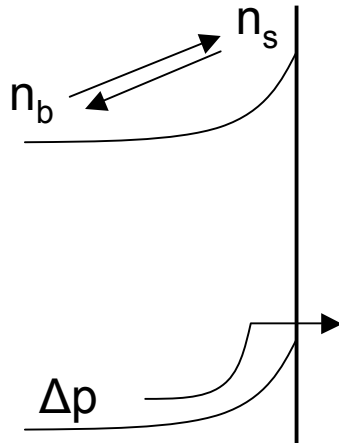
<Review>

Photocurrent: due to minority carrier current generated from photoinduced holes, Δp

$$h\nu \rightarrow \Delta n = \Delta p$$

$$n = N_D \gg \Delta n$$

$$\Delta p \gg p$$



If band bending is steep enough, all Δp cross the interface and J_{ph} is a constant

$$J = -J_0 [\exp(-qV / AkT) - 1] + J_{ph}$$

at Voc, $J = 0$ So,

$$J_0 [\exp(-qV / AkT) - 1] = J_{ph}$$

if $V > \frac{3kT}{q}$, we can ignore "1"

$$J_0 [\exp(-qV / AkT)] = J_{ph}$$

$$\therefore |V_{oc}| = \frac{AkT}{q} \ln \frac{J_{ph}}{J_0} \quad \Leftarrow \text{Name of this game is to minimize } J_0 \text{ to get large } V_{oc}$$

<Review>

SC/Metal interface,

$$J_0 = A^* T^2 \exp\left(-\frac{q\phi_b}{kT}\right) \quad : \text{ we want to have large } \phi_b \text{ to minimize } J_0$$

SC/Solution interface,

$$J_0 = qk_{et} [A] n_{s0}$$

$$J_{0,so\ln} < J_{0,m}$$

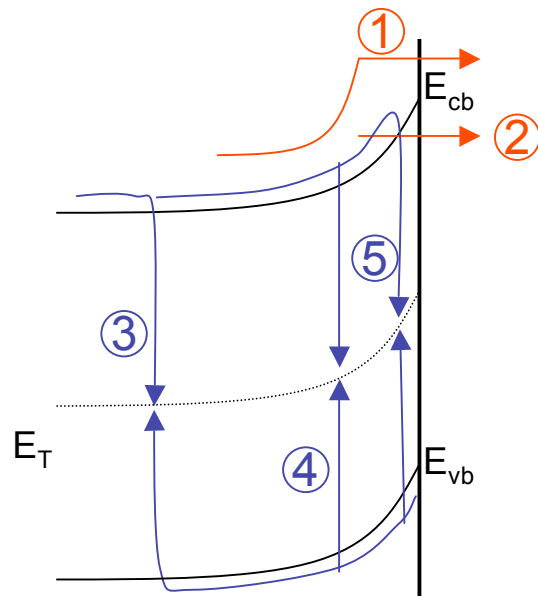
MIS,

insulator induces a less than unity probability for the e^- to cross the SC interface

Limit of ϕ_b ,

Max ϕ_b can only be as big as band gap energy, E_g

Five Mechanisms



Majority carrier process

① Thermionic Emission

② Tunneling

Minority carrier process

③ Bulk Recombination

④ Depletion Region Recombination

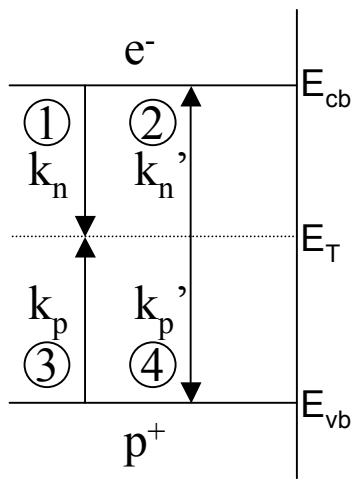
⑤ Surface Recombination

Under light or a bias, $np > n_i^2$. So there is a driving force to return to the equilibrium. Extra holes will recombine with extra electrons → Recombination

If somehow we were sucking carriers out of a region of the SC, $np < n_i^2$ and carriers would be generated → Generation

Shockley-Read-Hall statistics: Good model for all of these minority carrier processes of recombination in the SC

Bulk Recombination



$$r_n = \left. \frac{dn}{dt} \right|_R = \left. \frac{dn}{dt} \right|_1 + \left. \frac{dn}{dt} \right|_2 = -k_n \cdot n_b \cdot (1 - f_T) \cdot N_T + k_n' \cdot f_T \cdot N_T$$

$$r_p = \left. \frac{dp}{dt} \right|_R = \left. \frac{dp}{dt} \right|_3 + \left. \frac{dp}{dt} \right|_4 = -k_p \cdot p_b \cdot f_T \cdot N_T + k_p' \cdot (1 - f_T) \cdot N_T$$

Equilibrium conditions: Principle of detailed balance

➔ All fundamental processes and their inverse processes must self balance

$$\boxed{r_n = r_p = 0}$$

electrons,

$$k_n n_{b,0} (1 - f_T) N_T = k_n' f_T N_T$$

$$k_n' = \frac{1 - f_T}{f_T} k_n n_{b,0}$$

Bulk Recombination

For electrons,

$$k_n n_{b,0} (1 - f_T) N_T = k_n' f_T N_T$$

$$k_n' = \frac{1 - f_T}{f_T} k_n n_{b,0}$$

What is the ratio, $\frac{1 - f_T}{f_T}$?

Probability that a trap is filled is given by Fermi-Dirac Statistics.

$$f_T = \frac{1}{\exp[-(E_F - E_T)/kT] + 1}$$

$$1 - f_T = \frac{1}{\exp[(E_F - E_T)/kT] + 1}$$

$$\frac{f_T}{1 - f_T} = \frac{\frac{1}{\exp[-(E_F - E_T)/kT] + 1}}{1 - \frac{1}{\exp[-(E_F - E_T)/kT] + 1}} \times \frac{\exp[-(E_F - E_T)/kT] + 1}{\exp[-(E_F - E_T)/kT] + 1} = \frac{1}{\exp[-(E_F - E_T)/kT] + 1 - 1}$$

$$\therefore \frac{f_T}{1 - f_T} = \frac{1}{\exp[-(E_F - E_T)/kT]}$$

Bulk Recombination

$$\frac{1-f_T}{f_T} = \exp[-(E_F - E_T)/kT]$$

$$\therefore k_n' = k_n \left(\frac{1-f_T}{f_T} \right) n_{b,0} = k_n n_{b,0} \exp[-(E_F - E_T)/kT] = k_n N_C \exp[-(E_C - E_T)/kT]$$
$$(\because n_{b,0} = N_C \exp[-(E_C - E_F)/kT])$$

k_n' doesn't depend on the Fermi level and only depends on the trap energy level

For a given trap energy, $N_C \exp[-(E_C - E_T)/kT]$ is a constant.

Let $n_1 = N_C \exp[-(E_C - E_T)/kT]$

Then, $k_n' = k_n n_1$

Similarly, $k_p' = k_p p_1$ where $p_1 = N_V \exp[-(E_T - E_V)/kT]$

By plugging in k_n' and k_p' ,

$$r_n = \left. \frac{dn}{dt} \right|_R = -k_n \cdot n_b \cdot (1-f_T) \cdot N_T + k_n \cdot n_1 \cdot f_T \cdot N_T$$

$$r_p = \left. \frac{dp}{dt} \right|_R = -k_p \cdot p_b \cdot f_T \cdot N_T + k_p \cdot p_1 \cdot (1-f_T) \cdot N_T$$

Bulk Recombination

Steady state relationship: Major assumption in device problems and analysis

At steady state, $r_n = r_p$
(Different from the equilibrium)

$$U_{bulk} = -\left. \frac{dn}{dt} \right|_R = -\left. \frac{dp}{dt} \right|_R = \text{recombination rate}$$

Plug in our values for r_n & r_p

$$k_n \cdot n_b \cdot (1 - f_T) \cdot N_T - k_n \cdot n_1 \cdot f_T \cdot N_T = k_p \cdot p_b \cdot f_T \cdot N_T - k_p \cdot p_1 \cdot (1 - f_T) \cdot N_T$$

N_T s cancel and solve for f_T

$$\rightarrow k_n \cdot n_b \cdot (1 - f_T) - k_n \cdot n_1 \cdot f_T = k_p \cdot p_b \cdot f_T - k_p \cdot p_1 \cdot (1 - f_T)$$

$$\rightarrow k_n \cdot n_b - k_n \cdot n_b \cdot f_T - k_n \cdot n_1 \cdot f_T = k_p \cdot p_b \cdot f_T - k_p \cdot p_1 + k_p \cdot p_1 \cdot f_T$$

$$f_T = \frac{k_n \cdot n_b + k_p \cdot p_1}{k_n \cdot n_b + k_n \cdot n_1 + k_p \cdot p_b + k_p \cdot p_1}$$

$$1 - f_T = \frac{k_p \cdot p_b + k_n \cdot n_1}{k_n \cdot n_b + k_n \cdot n_1 + k_p \cdot p_b + k_p \cdot p_1}$$

Now that we have f_T and $1 - f_T$ we can use them to calculate U_{bulk} .

Bulk Recombination

$$\begin{aligned}U_{bulk} &= -\left.\frac{dn}{dt}\right|_R = -\left.\frac{dp}{dt}\right|_R \\&= -k_n \cdot n_b \cdot (1 - f_T) \cdot N_T + k_n \cdot n_1 \cdot f_T \cdot N_T \\U_{bulk} &= N_T \frac{k_n \cdot n_b \cdot (k_p \cdot p_b + k_n \cdot n_1) - k_n \cdot n_1 \cdot (k_n \cdot n_b + k_p \cdot p_1)}{k_n \cdot (n_b + n_1) + k_p \cdot (p_b + p_1)} \\U_{bulk} &= N_T \frac{k_n \cdot k_p \cdot (n_b \cdot p_b - n_1 \cdot p_1)}{k_n \cdot (n_b + n_1) + k_p \cdot (p_b + p_1)} \\n_1 \cdot p_1 &= N_C \exp[-(E_C - E_T)/kT] \cdot N_V \exp[-(E_T - E_V)/kT] \\&= N_C \cdot N_V \cdot \exp[-(E_C - E_V)/kT] = n_i^2\end{aligned}$$

So the recombination rate in bulk,

$$U_{bulk} = N_T \frac{k_n \cdot k_p \cdot (n_b \cdot p_b - n_i^2)}{k_n \cdot (n_b + n_1) + k_p \cdot (p_b + p_1)}$$